Experiment 9: AC circuits

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INTRO TO EXPERIMENTAL PHYS-LAB 1493/1494/2699

Introduction

- Last week (RC circuit):
 - **<u>Constant</u>** Voltage power source (constant over time)
- This week:
 - A new component: the inductor
 - Alternating Current (AC) circuits
 - Time dependent voltage source
 - Leads to:
 - Time dependent currents (alternating currents)
 - Phase shifts in voltage and currents in components with respect to one another
 - Resonance

Varying electromagnetic fields

• The 2nd and 4th Maxwell equations in vacuum and with no sources read: $\vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t}$

$$\vec{\nabla} \times \vec{B} = \frac{1}{c^2} \frac{\partial \vec{E}}{\partial t}$$

- Varying magnetic field generates electric field and viceversa!
- We therefore expect from a changing B field to give rise to a current or potential difference
- This is summarized by **Faraday's law**:



$$\mathcal{E} = -\frac{d\Phi_B}{dt}$$

- Change in magnetic flux creates an e.m.f.
- Magnetic flux can be varied in two ways:
 - 1. Change the B field
 - 2. Change the surface

Introducing the inductor

- It <u>stores energy</u> in form of magnetic fields (analogous to the capacitor)
- From Faraday's law one deduces the expression for the potential difference at the two ends of the inductor:

$$V_L = -L\frac{dI}{dt}$$

- The inductor is only sensitive to the change in current! No change = no voltage
- Negative sign indicates that the inductor opposes any change in current (Lenz's Law)





Why AC circuits?

- Sensitive to input frequency (*i.e.* function generator frequency)
- Serve as <u>signal frequency</u> <u>filters</u>:
 - High-frequency filters
 - Low-frequency filters
 - Band-pass filters
- Transformers
 - <u>Induction effects</u> Ability to raise or lower the voltage amplitude.
- Generators and Motors







e.g. Speakers

e.g. Radios

AC circuits: resistors

- AC circuits have an enormous range of applications. Here we cover the most important aspects
- For this lab:
 - Consider only sources that vary sinusoidally:

$$I(t) = I_{\max}\sin(\omega t)$$

- Simple example:
 - Function generator + resistor
- **Ohm's Law:** voltage across the resistor is just $V_R = IR$



AC circuits: capacitors

- More interesting case: connect a capacitor to the AC voltage source
- Last time we saw that the voltage across a capacitor is given by:

• Therefore, when the current is sinusoidal the voltage is given by:

 $V_C(t)$

I(t)

 $=\frac{1}{C}\int I_{\max}\sin(\omega t)dt$

 $= -\frac{I_{\max}}{\omega C}\cos(\omega t)$

 $V_C(t) = \frac{1}{C} \int I(t) dt$



PHYS 1493/1494/2699: Exp. 9 – AC circuits

AC circuits: inductors

• The voltage is still sinusoidal:

$$V_L(t) = L \frac{dI}{dt}$$

= $\omega L I_{\max} \cos(\omega t)$
= $V_{\max} \sin\left(\omega t + \frac{\pi}{2}\right)$

- Inductor voltage is also <u>phase</u> <u>shifted</u> w.r.t. current.
- Voltage across the inductor anticipates the current through it by: $\phi = \pi/2$



Voltage maxima: a closer look

• Given our expression for V_R , the *maximum value* of the voltage across the <u>resistor</u> is just given by Ohm's Law:

$$V_R^{\max} = I_{\max}R$$

• The maximum voltage across the <u>capacitor</u> is a function of ω :

$$V_C^{\max} = \frac{1}{\omega C} I_{\max} \equiv I_{\max} X_C$$

Capacitive reactance



Given an oscillating input current the capacitor voltage is higher for small frequencies and lower for high frequencies

 But, the maximum voltage across the inductor is also a function of the driving frequency:

$$V_L^{\max} = \omega L I_{\max} \equiv I_{\max} X_L$$

Inductive reactance

$$X_L = \omega L$$

The inductor voltage is instead higher for large frequencies and lower for small ones

Physical explanation: capacitors

- <u>Question</u>: Why does the capacitor resist low-frequency signals more than high-frequency ones?
- Last time: when charging/discharging the capacitor, the current the rate at which you can charge it – decreases exponentially. It becomes harder and harder to push in more charge as the capacitor fills up.



Physical explanation: capacitors



- Rapidly varying signals (high frequency) quickly charge/discharge capacitor before it fills with charge → low impedance.
- Slowly varying signals (low frequency) charge the capacitor to its limit, slowing down the rate: that is, decreasing the current!
- Now that we have introduced the language of reactances, <u>you can</u> <u>think about the capacitor somehow as a resistor with ω-dependent</u> <u>resistance</u>

Physical explanation: inductors

- Question: Why does the inductor resist high-frequency signals more than low-frequency ones?
- Think about the nature of an inductor: it is a *coil of wire*. If the current in the wire changes, then the magnetic flux through the coil changes \rightarrow induction!
- <u>Lenz's Law:</u> a coil will oppose changes in magnetic flux. Self-induced EMF is:

$$\mathcal{E} = \bigcirc N \frac{d\Phi_B}{dt} = \bigcirc L \frac{dI}{dt}$$



A life spent after a minus sign... and everyone always forgets about it!

(Heinrich Lenz)

- Rapidly varying signals strongly change the flux, so the inductor "pushes back" harder against the flow of current!
 - Voltage is maximum (and opposing) when I changes most rapidly (high frequency)
 - Voltage = 0 when I is constant (low frequency)

RLC circuits

 Let's see what happens when we combine all these three components in a series:



From <u>Kirchhoff's first law</u> (loops):

Use what we learned about inductors and capacitors a fe

slides ago

$$V(t) = V_R + V_L + V_C = IR + L\frac{dI}{dt} + \frac{Q}{C}$$

$$= I_{\max} \left[R\sin\omega t + X_L\sin\left(\omega t + \frac{\pi}{2}\right) + X_C\sin\left(\omega t - \frac{\pi}{2}\right) \right]$$

$$= I_{\max} \left[R\sin\omega t + (X_L - X_C)\cos\omega t \right] \quad (\omega t \pm \pi/2) = \pm \cos(\omega t)$$

$$= I_{\max} \left[R\sin\omega t + (X_L - X_C)\cos\omega t \right] \quad (\omega t \pm \pi/2) = \pm \cos(\omega t)$$

RLC circuits: phase shift

- After passing through the three components the voltage will have some phase shift
- Let's then impose to V(t) to look like:

 $V(t) = V_{\max} \sin(\omega t + \overline{\phi})$



 $=V_{\max}\left[\sin(\omega t)\cos\phi + \cos(\omega t)\sin\phi\right]$

- <u>Comparing</u> with the equation from the previous slide it must necessarily be: $\begin{cases} V_{\max} \cos \phi = I_{\max} R \\ V_{\max} \sin \phi = I_{\max} (X_L - X_C) \end{cases}$

And hence the phase shift is: $\tan \phi = \frac{X_L - X_C}{R} = \frac{\omega L - 1/\omega C}{R}$

The phase shift will depend both on the characteristics of the circuit (R, C, L) and on the **frequency** of the input signal!

RLC circuits: phase shift

- What about the maximum amplitude for the voltage?
- Let take again:

$$V_{\max} \cos \phi = I_{\max} R$$
$$V_{\max} \sin \phi = I_{\max} (X_L - X_C)$$



Let's now square both equations and add them together:

$$V_{\max} = I_{\max} \sqrt{R^2 + (X_L - X_C)^2}$$
$$= I_{\max} Z$$

- The quantity Z is called the impedance of the RLC circuit
- <u>NOTE</u>: the previous equation resembles very closely Ohm's law for resistors!
- This procedure can actually be generalized introducing the socalled <u>phasor formalism</u>

Resonant frequency

- So the whole RLC system has this peculiar frequency dependent "effective resistance". In particular:
 - <u>High-frequencies</u>: killed by the *inductor*
 - Low-frequencies: killed by the capacitor
- We therefore expect to have a particular frequency (ω_0) in the middle range that goes through the system almost untouched $I_{\max} = \frac{V_{\max}}{Z} = \frac{V_{\max}}{\sqrt{R^2 + (X_L X_C)^2}}$



High Resonant

- For a given input voltage, the current in the circuit is maximum when Z is minimum *i.e.* when $X_L = X_C$.
- The *resonant frequency* is given by:

$$\omega_0 L = \frac{1}{\omega_0 C} \implies \omega_0 = \frac{1}{\sqrt{LC}}$$

Resonance



<u>Terminology</u>

FWHM: **<u>Full Width</u>** at <u>**Half Maximum**</u> *Is the full width of the resonance peak at the point where its height is halfway between zero and the maximum.*

Recall that the resistor voltage V_R is directly proportional to the magnitude of current.

The Experiment

Main goals

Resonance of RLC circuit:

 Measure the resonant frequencies and FWHM for three known circuits



- Compute the unknown inductance of a copper coil by finding the resonant frequency of the whole system
- Observe the phase shift, φ, between the driving signal and the three components (R, L and C) of the circuit
- Compare with expected value

Experimental setup



Experimental setup

- Recommendations:
 - Set the function generator peak-to-peak voltage to 20 V, the maximum allowed.
 - There is a 0-2V / 0-20V selector button in addition to the voltage knob.
 - Make sure the <u>oscilloscope is set to trigger on channel 1</u>, the function generator signal. You can do this by pressing the TRIGGER button and checking in the window menu that CH 1 is selected.
 - Use the MEASURE tools to observe peak-peak amplitudes, signal periods, and signal frequencies. Let the scope do the work for you!
 - Make sure that both peaks are in the viewable range of the scope!

Resonance measurements

- Set the oscilloscope to look at the potential between the two ends of the resistor
- First localize ω_0 by looking at when the amplitude of the signal gets amplified
- Then, for about 20 frequencies above and below ω_0 , record peak-peak voltage across resistor
- Normalize your values such that the maximum is $V_{pp} = 1$
- Repeat for three values of the resistance (10 Ω, 50 Ω, 500 Ω)
- Plot V_{pp} vs. ω as shown in the figure



RLC Resonance

Resonance measurements

- Use the plots to determine the value of the resonant frequencies (with errors!)
- Also use the plots to measure the FWHM of the three curves.
- <u>Compare the results from the</u> <u>three measurements</u>
- When you're finished with this part, <u>replace the inductor with</u> <u>the large copper coil</u>
- Repeat the previous measures and from the value of the average resonant frequency compute the inductance, *L*, of the wire



Phase shift measurement

- Replace the copper ring with the known capacitor again and set $R = 30 \Omega$
- Locate the resonant frequency and for 5 values at, above and below it measure the phase shift across the resistor
- Now increase the frequency well above the resonant one. This makes the inductor much more important than the capacitor. Measure φ
- Now make the capacitor more important by going way below the resonant frequency. Measure φ again



$$\phi = 2\pi \frac{\iota_R - \iota_d}{T_d}$$

Compare the results obtained with the expected ones

Tips

- Don't get confused! The frequency reported by the oscilloscope (in the MEASURE mode) is *f*. It is related to what we called frequency so far by $\omega = 2\pi f$.
- Once you manage to locate the maximum of the V_{pp} curve and you took 20 points above and below it, try to take more points right around you maximum. This will reduce your uncertainty on ω_0
- When plotting V_{pp} vs ω remember to take enough data to measure the FWHM. When R is large, the peak is very broad. Keep taking data until you pass half height
- The resistor is old but if you look carefully on each knob there is a label telling you how many Ω correspond to that knob.