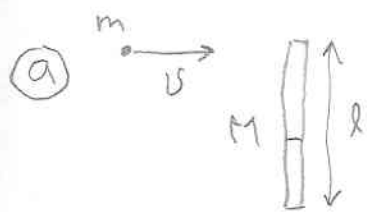


2. Work out two examples

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mass m moving with velocity $v \perp$ to rod of length l + mass M (at rest)

strikes and sticks to the end of the rod. Describe the resulting motion.

* CM motion moves with $(m+M)\vec{V}_{cm} = m\vec{v}$
or $\vec{V}_{cm} = \frac{m}{m+M} \vec{v}$

Note location of cm must be d above center where $(m+M)d = m\frac{l}{2} + M \cdot 0 \Rightarrow d = \frac{m}{m+M} \frac{l}{2}$

* Rotation about CM

$$L_{int} = m v \left(\frac{l}{2} - d \right) = \frac{m M v}{m+M} \frac{l}{2}$$

$$L_{final} = I_{cm} \omega$$

$$I_{cm} = \underbrace{m \left(\frac{l}{2} - d \right)^2}_{I \text{ for } m} + \underbrace{\frac{1}{12} M l^2 + M d^2}_{\text{parallel axis theorem for rod about new CM}}$$

parallel axis theorem for rod about new CM

December 01, 2020

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$$I_{cm} = \left\{ \frac{m M^2}{(m+M)^2} + \frac{1}{3} M + \frac{M m^2}{(m+M)^2} \right\} \frac{l^2}{4}$$

$$= \frac{3 M m + M(m+M)}{m+M} \frac{l^2}{12} = \frac{(4m+M)M}{m+M} \frac{l^2}{12}$$

$$\omega = \frac{L_{int}}{I_{cm}} = \frac{m+M}{(4m+M)M} \frac{12}{l^2} \times \frac{m M v}{m+M} \frac{l}{2}$$

$$= \frac{6m}{4m+M} \frac{v}{l}$$



A bowling ball of mass M and radius R is released moving at velocity v_0 and begins to slide with coefficient of friction μ . What is its final velocity

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$$M \frac{dv}{dt} = -\mu M g \quad v(t) = v_0 - \mu g t$$

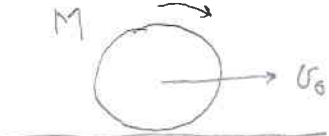
$$I \dot{\omega} = \tau = \mu M g R$$

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(b)  A bowling ball of mass M and radius R is released moving at velocity v_0 and begins to slide with coefficient of friction μ . What is its final velocity

$$M \frac{dv}{dt} = -\mu Mg \quad v(t) = v_0 - \mu g t$$

$$I \dot{\omega} = \tau = \mu Mg R$$

$$\dot{\omega} = \frac{\mu Mg R}{I} = \frac{\mu Mg R}{\frac{2}{5} M R^2} = \frac{5\mu g}{2R}$$

$$\omega(t) = \frac{5\mu g}{2R} t$$

stops sliding when

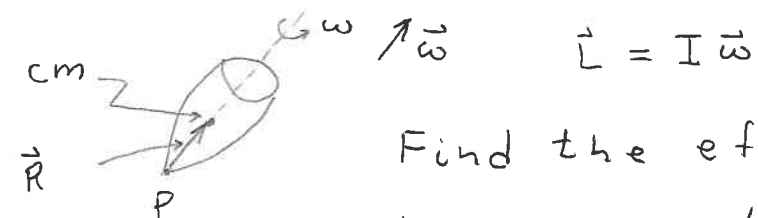
$$R\omega(t) = v(t)$$

$$\frac{5}{2} \mu g t = v_0 - \mu g t$$

$$\frac{7}{2} \mu g t = v_0, \quad t = \frac{2}{7} \frac{v_0}{\mu g}$$

$$v = \omega R = \frac{5\mu g}{2} \frac{2}{7} \frac{v_0}{\mu g} = \frac{5}{7} v_0$$

3. Motion of rapidly spinning top



mass M

Find the effect of the torque exerted by gravity about P . Assume $\vec{\omega}$ is

so large that any small additional rotation of the top can be neglected.

$$\dot{\omega} = \frac{\mu M g R}{I} = \frac{\mu M g R}{\frac{2}{5} M R^2} = \frac{5 \mu g}{2 R}$$

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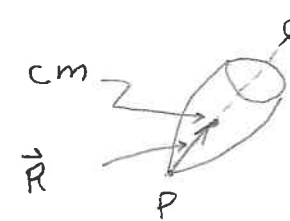
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3. Motion of rapidly spinning top



$$\vec{L} = I \vec{\omega}$$

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$$\frac{d\vec{L}}{dt} = \vec{\tau} = \vec{R} \times (-mg \hat{z})$$

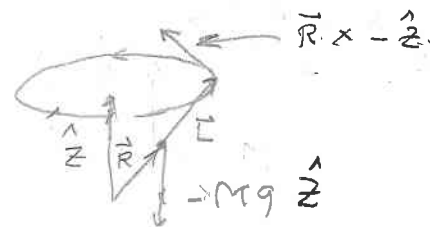
$$\text{Use } \vec{R} = I \vec{\omega} \frac{R}{I \omega} = \vec{L} \frac{R}{I \omega}$$

$$\frac{d\vec{L}}{dt} = \frac{R}{I \omega} \vec{L} \times (-Mg \hat{z}) = \frac{MgR}{I \omega} \hat{z} \times \vec{L}$$

$\vec{\omega}_{\text{prec}}$

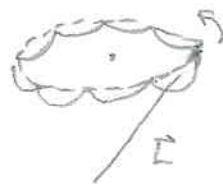
$\Rightarrow \vec{L}$ precesses with angular velocity

$$\vec{\omega}_{\text{prec}} = \frac{MgR}{I \omega} \hat{z}$$



nutation

Note if rotation axis is held fixed and then released \vec{L} will begin to rotate but dip down:



Nutation - added angular momentum associate with precession, $\Delta L \propto +\hat{z}$ comes

from decrease in L_z from spin on original axis.

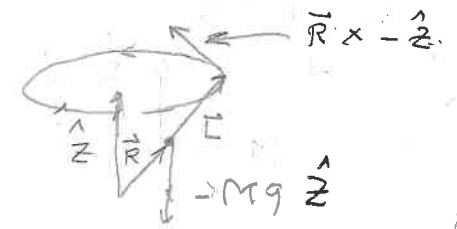
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Use $\vec{R} = I\vec{\omega} \frac{R}{I\omega} = \vec{L} \frac{R}{I\omega}$

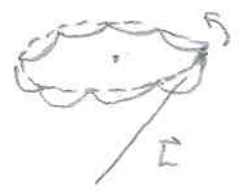
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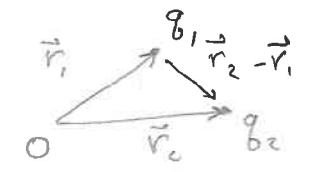


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IV Electrostatics

A. Coulomb's law



1. Introduction:

Inverse square law force between to charges q_1 and q_2

$$\vec{F}_{ion2} = k \frac{\vec{r}_2 - \vec{r}_1}{|\vec{r}_2 - \vec{r}_1|^2} \frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|^2}$$

Two common sets of units

a) esu $k = 1$ $1 \text{ esu} = \sqrt{\text{dyne cm}}$

b) SI: $k = 8.9875 \times 10^9 \frac{\text{Newton} \cdot \text{m}^2}{(\text{Coulomb})^2}$

esu & Coulomb are unit of charge

Note force between two 1 Coulomb charges 1 cm apart is

$$\approx \left[9 \times 10^9 \text{ Newton} \cdot \text{m}^2 \right] \times \frac{1}{(0.01 \text{ m})^2} \times \frac{10^5 \text{ dyne}}{\text{Newton}}$$

$$= 9 \times 10^{18} \text{ dyne}$$

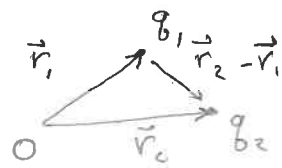
$$1 \text{ Coulomb} \equiv 3 \times 10^9 \text{ esu}$$

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(166)

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(167)

Like gravity, the Coulomb force is a fundamental force that comes from forces between elementary particles at the atomic level. Because, like gravity it is long range, $F \propto \frac{1}{r^2}$, it is easily recognized at a macroscopic scale

- the strong nuclear force $F_{\text{strong}} \sim \frac{e^{-\frac{r}{10^{-13} \text{ cm}}}}{r^2}$

- the weak force $F_{\text{weak}} \sim \frac{e^{-\frac{r}{10^{-16} \text{ cm}}}}{r^2}$

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2. Principle of superposition:

Just as with gravity, electrostatic forces add $\vec{F}_{1+2 \text{ on } 3} = \vec{F}_{1 \text{ on } 3} + \vec{F}_{2 \text{ on } 3}$

Given N charges q_1, q_2, \dots, q_N at positions $\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N$ their force on an $N+1^{\text{st}}$ charge q at \vec{r} is given by

$$\vec{F}_{\text{on } q} = \sum_{i=1}^N \frac{q_i q}{|\vec{r}_q - \vec{r}_i|^3} (\vec{r}_q - \vec{r}_i)$$

[This simple linear superposition fails for extremely strong fields near point-like atomic charges.]

3. Define the electric field $\vec{E}(\vec{r})$:

Since $\vec{F}_{\text{on } q}$ is proportional to q

we can factor out q and define

$$\vec{E}(\vec{r}) = \lim_{q \rightarrow 0} \frac{1}{q} \vec{F}_{\text{on } q} = \sum_{i=1}^N \frac{\vec{r} - \vec{r}_i}{|\vec{r} - \vec{r}_i|^3} q_i$$

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The limit $\delta q \rightarrow 0$ allows for the possibility that δq might shift the position of the other charges. $\vec{E}(\vec{r})$ does not depend on the small test charge used to measure it.

While $\vec{E}(\vec{r})$ looks like a foolish rescaling of $\vec{F}_{\text{on } q}$ its introduction is profound!

- The potential energy of the assembled N charges

$$E = \frac{1}{2} \sum_{i \neq j} \frac{q_i q_j}{|\vec{r}_i - \vec{r}_j|} = \int d^3r \underbrace{\frac{\vec{E}(\vec{r})^2}{8\pi}}_{\text{energy density!}}$$

- Introduction of $\vec{E}(\vec{r})$ can make electrostatic become causal. modify relation between \vec{r}_i & \vec{E} to allow changes in \vec{r}_i to propagate to give changes in $\vec{E}(\vec{r})$

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4. Recall result for potential energy

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for gravity $U(\vec{r}_1, \vec{r}_2) = -\frac{Gm_1 m_2}{|\vec{r}_1 - \vec{r}_2|}$

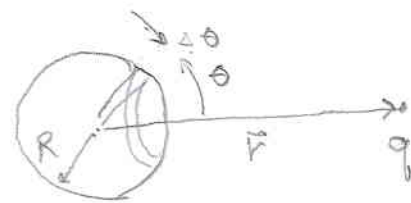
$$\vec{F}_{1 \text{ on } 2} = -\vec{\nabla}_{\vec{r}_2} \left[-\frac{Gm_1 m_2}{\sqrt{(x_1 - x_2)^2 + (y_1 - y_2)^2 + (z_1 - z_2)^2}} \right]$$

$$= -\frac{Gm_1 m_2 (\vec{r}_2 - \vec{r}_1)}{|\vec{r}_1 - \vec{r}_2|^3} \quad \checkmark$$

for two charges $U(r_1, r_2) = +\frac{q_1 q_2}{|\vec{r}_1 - \vec{r}_2|}$

for N charges $U(\vec{r}_1, \dots, \vec{r}_N) = \frac{1}{2} \sum_{i \neq j} \frac{q_i q_j}{|\vec{r}_i - \vec{r}_j|}$

5. Find the potential energy of a charge q displaced by \vec{r} from the center of a sphere (not a ball) of charge Q with radius R



$$|\vec{r} - \vec{r}'|^2 = r^2 + R^2 - 2rR \cos \theta$$

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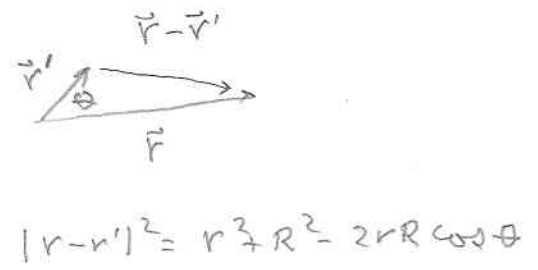
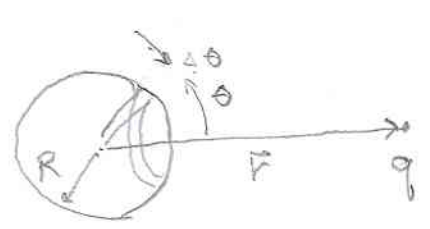
$$\vec{F}_{q \text{ on } 2} = -\vec{\nabla}_{\vec{r}_2} \left[-\frac{Gm_1m_2}{\sqrt{(x_1-x_2)^2 + (y_1-y_2)^2 + (z_1-z_2)^2}} \right]$$

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$$|r-r'|^2 = r^2 + R^2 - 2rR \cos \theta$$

Add up the potential energy of each ring of charge, equidistant from q :

$$U(r) = \int_0^\pi \underbrace{d\theta R \cdot 2\pi R \sin \theta}_{\text{area of ring}} \underbrace{\frac{Q}{4\pi R^2}}_{\text{surface density of charge on sphere}} \frac{q}{\sqrt{R^2 + r^2 - 2Rr \cos \theta}}$$

Let $z = \cos \theta$

$$= \frac{Qq}{2} \int_{-1}^1 dz \frac{1}{\sqrt{R^2 + r^2 - 2Rrz}}$$

$$= \frac{Qq}{2} \left(-\frac{1}{2Rz} \right) 2 \left\{ \frac{\sqrt{R^2 + r^2 - 2Rr}}{|R-r|} - \frac{\sqrt{R^2 + r^2 + 2Rr}}{|R+r|} \right\}$$

- 2x smaller (R,r)

$$= Qq \frac{1}{\text{larger}(R,r)}$$

