

$$\vec{B} = -\frac{I}{c} \oint_C \frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} \times d\vec{l}$$

$$\vec{B} = \hat{z} \frac{I}{c} \frac{1}{z^2+R^2} \times \frac{R}{\sqrt{z^2+R^2}} \times 2\pi R$$

$$= \frac{2\pi I}{c} \frac{R^2}{(z^2+R^2)^{3/2}} \hat{z}$$

4. Multipole expansion

The behavior of $\vec{E}(\vec{r})$ & $\vec{B}(\vec{r})$ far from localized charges and currents which produces the is simple to describe

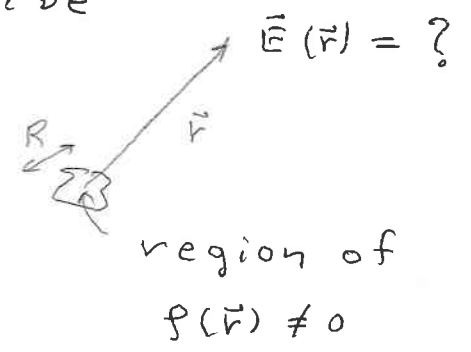
a) Easiest for $\vec{E}(\vec{r})$

$$\vec{E}(\vec{r}) = \int d^3r' \rho(r') \frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3}$$

for r large, $r \gg R$

$$\approx \frac{\vec{r}}{r^3} \int d^3r' \rho(r')$$

$$= Q \frac{\vec{r}}{r^3}$$



Looks like a point charge at large distance

Find next term in r/R expansion

$$\frac{1}{|\vec{r}-\vec{r}'|^3} = \frac{1}{[r^2 - 2\vec{r}\cdot\vec{r}' + r'^2]^{3/2}} \approx \frac{1}{r^3} \left[1 - \frac{2\vec{r}\cdot\vec{r}'}{r^2} \right]^{3/2}$$

$$\approx \frac{1}{r^3} \left[1 + \frac{3}{2} \frac{2\vec{r}\cdot\vec{r}'}{r^2} + \dots \right]$$

$$\Rightarrow \vec{E}(\vec{r}) = \int d^3r' \frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} \approx \int d^3r' \frac{\vec{r}-\vec{r}'}{r^3} \left[1 + 3 \frac{\vec{r}\cdot\vec{r}'}{r^2} \right] \rho(r')$$

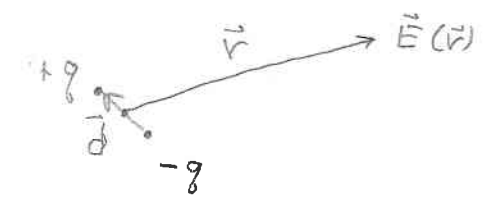
$$= \int d^3r' \rho(r') \left\{ \frac{\vec{r}}{r^3} + \vec{r} \frac{3\vec{r}\cdot\vec{r}'}{r^5} - \frac{\vec{r}'}{r^3} + \dots \right\}$$

Use $Q = \int d^3r' \rho(r')$
charge

$\vec{D} = \int d^3r' \rho(r') \vec{r}'$
electric dipole moment

$$\vec{E}(r) = \underbrace{Q \frac{\vec{r}}{r^3}}_{\text{monopole}} + \underbrace{\frac{3\hat{r}(\hat{r}\cdot\vec{D}) - \vec{D}}{r^3}}_{\text{dipole}} + \text{quadrupole}$$

Easy example of a dipole



Find next term in r/R expansion

$$\frac{1}{|\vec{r}-\vec{r}'|^3} = \frac{1}{[\vec{r}^2 - 2\vec{r}\cdot\vec{r}' + \vec{r}'^2]^{3/2}} \approx \frac{1}{r^3} \left[1 - \frac{2\vec{r}\cdot\vec{r}'}{r^2} \right]^{3/2}$$

$$\approx \frac{1}{r^3} \left[1 + \frac{3}{2} \frac{2\vec{r}\cdot\vec{r}'}{r^2} + \dots \right]$$

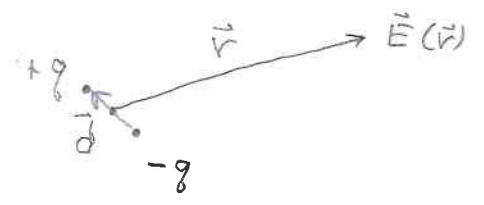
$$\Rightarrow \vec{E}(\vec{r}) = \int d^3r' \frac{\vec{r}-\vec{r}'}{|\vec{r}-\vec{r}'|^3} \rho(\vec{r}') \approx \int d^3r' \frac{\vec{r}-\vec{r}'}{r^3} \left[1 + 3 \frac{\vec{r}\cdot\vec{r}'}{r^2} \right] \rho(\vec{r}')$$

$$= \int d^3r' \rho(\vec{r}') \left\{ \frac{\vec{r}}{r^3} + \vec{r}' \frac{3\vec{r}\cdot\vec{r}'}{r^5} - \frac{\vec{r}'}{r^3} + \dots \right\}$$

Use $Q = \int d^3r' \rho(\vec{r}')$ charge
 $\vec{D} = \int d^3r' \rho(\vec{r}') \vec{r}'$ electric dipole moment

$$\vec{E}(\vec{r}) = \underbrace{Q \frac{\vec{r}}{r^3}}_{\text{monopole}} + \underbrace{\frac{3\hat{r}(\hat{r}\cdot\vec{D}) - \vec{D}}{r^3}}_{\text{dipole}} + \text{quadrupole}$$

Easy example of a dipole



$$E(\vec{r}) = q \frac{\vec{r} - \frac{d}{2}\hat{z}}{|\vec{r} - \frac{d}{2}\hat{z}|^3} - q \frac{\vec{r} + \frac{d}{2}\hat{z}}{|\vec{r} + \frac{d}{2}\hat{z}|^3}$$

$$\approx q \left(\vec{r} - \frac{d}{2}\hat{z} + 3 \frac{\vec{r} \cdot \frac{d}{2}\hat{z}}{r^3} \right) - q \left(\vec{r} + \frac{d}{2}\hat{z} - 3 \frac{\vec{r} \cdot \frac{d}{2}\hat{z}}{r^3} \right)$$

$$= q \frac{1}{r^3} [3\hat{r}(\hat{r}\cdot\vec{d}) - \vec{d}]$$

\Rightarrow dipole moment $\vec{D} = q\vec{d}$

b) Since \vec{E} & \vec{B} both obey the same equations when $\rho(\vec{r})=0$ & $\vec{j}(\vec{r})=0$ the same expansion must hold for $\vec{B}(\vec{r})$ for large r

$$\vec{B}(\vec{r}) = q^* \frac{\vec{r}}{r^3} + \frac{3\hat{r}(\hat{r}\cdot\vec{\mu}) - \vec{\mu}}{r^3}$$

where $\vec{\mu}$ is the magnetic dipole moment

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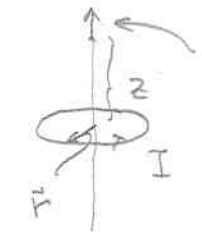
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where $\vec{\mu}$ is the magnetic dipole moment

Find $\vec{\mu}$ for our loop of current



$$B(z) = \hat{z} \frac{2\pi I R^2}{c r^3}$$

$$= \frac{3 \vec{\mu} \cdot \hat{r} \hat{r} - \vec{\mu}}{r^3} \quad | \quad \vec{r} = r \hat{z}$$

symmetry requires $\vec{\mu} = \mu \hat{z}$

$$\frac{2\pi I R^2}{c r^3} = \frac{2\mu}{r^3} \Rightarrow \mu = \frac{\pi R^2 I}{c}$$

Example: A mass M with axial symmetry has uniform mass and charge density & rotates about its axis of symmetry with angular velocity $\vec{\omega}$, then $\vec{\mu}$ & \vec{L} are proportional.

First consider our ring:

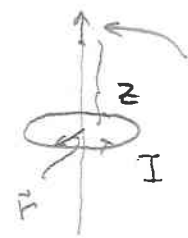
$$M = \frac{\pi R^2 I}{c} \quad L = MR^2 \omega$$

$$I = \frac{Q}{2\pi \text{one rotation}} = \frac{Q}{2\pi/\omega}$$

gyromagnetic ratio

$$\mu = \frac{\pi R^2}{c} \frac{Q}{2\pi} \omega = \frac{Q}{2Mc} (MR^2 \omega) = \frac{Q}{2Mc} L$$

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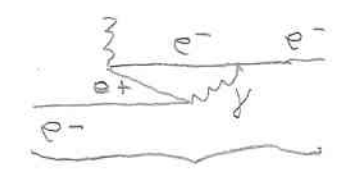
$$\mu = \frac{\pi R^2}{c} \frac{Q}{2\pi} \omega = \frac{Q}{2Mc} (MR^2 \omega) = \frac{Q}{2Mc} L$$

This will also be true for our axially symmetric solid with proportional mass and charge density - divide it into rings!

For an electron

$$\vec{\mu}_{e^-} = \left[2 + \frac{d}{\pi} + \dots \right] \frac{e}{2m_{e^-} c} \vec{L}$$

$$d \approx \frac{1}{137}$$



mixture of e^- & $2e^- + e^+$

$$\mu_{\text{mass}} \neq \mu_{\text{charge}}$$

For a proton

$$\vec{\mu}_p = 5.583 \frac{e}{2m_p c} \vec{L} \quad \mu_{\text{mass}} \neq \mu_{\text{charge}}$$



$$q_u = \frac{2}{3} e$$

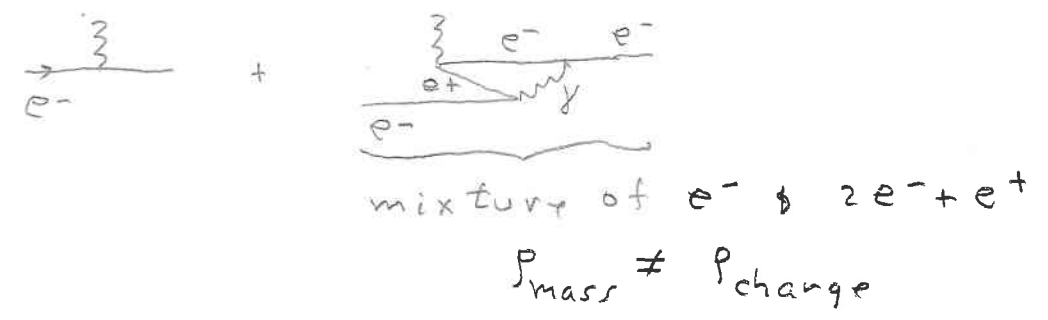
$$q_d = -\frac{1}{3} e$$

This will also be true for our axially symmetric solid with proportional mass and charge density - divide it into rings!

For an electron

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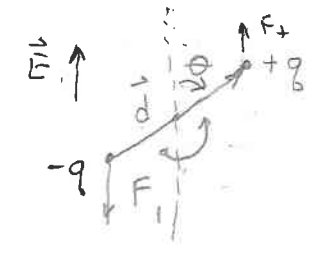


For a proton

$$\vec{\mu}_p = 5.583 \frac{e}{2m_p c} \vec{L} \quad P_{mass} \neq P_{charge}$$



Finally ask what forces act on a dipole if it is placed in a uniform field



It experiences a torque $|\tau| = \left[\frac{d}{2} \sin \theta E q \right] \times 2$

$$\text{or } \vec{\tau} = \vec{D} \times \vec{E}$$

Since \vec{E} , \vec{D} and \vec{B} , $\vec{\mu}$ are in close analogy

$$\vec{B} \uparrow \quad \vec{\mu} \quad \vec{\tau} = \vec{\mu} \times \vec{B}$$

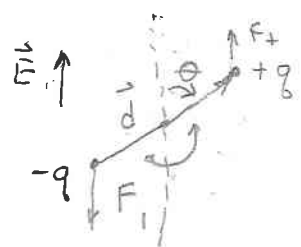
If $\vec{\mu}$ comes from rotation and angular momentum, then

$$\vec{\mu} = \frac{Q}{2mc} \vec{L} \quad \vec{\tau} = \vec{\mu} \times \vec{B}$$

$$\text{or } \frac{d\vec{L}}{dt} = \vec{\tau} = \vec{\mu} \times \vec{B} = \frac{Q}{2mc} \vec{L} \times \vec{B}$$

$$\text{or } \frac{d\vec{L}}{dt} = - \left[\frac{QB}{2mc} \right] \times \vec{L} \Rightarrow \vec{\omega}_{prec} = \frac{QB}{2mc}$$

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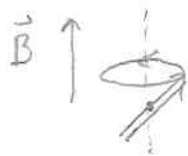
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$\vec{\mu} = \frac{q}{2mc} \vec{L}$

Larmor precession $\omega_{prec} = \frac{qB}{2mc}$

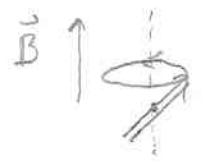
Can impose a radio frequency E & M field with frequency ω & see energy absorbed when $\omega = \omega_{prec}$.

Since most of the atomic nuclei in your body each have $\vec{\mu} \neq 0$ a spatially varying \vec{B} field will lose energy only to those nuclei at \vec{r} with $\frac{qB(\vec{r})}{2mc} = \omega_{RF}$



Magnetic resonant MRI imaging

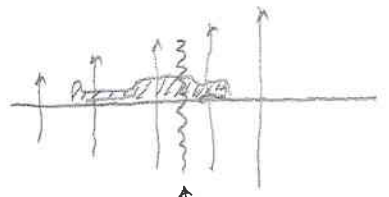
use gradients in 3 dimensions and do a lot of image processing



Larmor precession

$$\vec{\mu} = \frac{Q}{2mc} \vec{L} \quad \omega_{\text{prec}} = \frac{QB}{2mc}$$

Can impose a radio frequency E & M field with frequency ω & see energy absorbed when $\omega = \omega_{\text{prec}}$. Since most of the atomic nuclei in your body each have $\vec{\mu} \neq 0$ a spatially varying \vec{B} field will lose energy only to those nuclei at \vec{r} with $\frac{QB(\vec{r})}{2mc} = \omega_{\text{RF}}$



magnetic resonant MRI imaging

$$\omega_{\text{prec}} = \frac{QB(\vec{r})}{2mc}$$

use gradients in 3 dimensions and do a lot of image processing

5. How does a particle with charge q move in a uniform magnetic field \vec{B} ?

$$\frac{d\vec{p}}{dt} = \frac{q}{c} \vec{v} \times \vec{B}$$

$$= \frac{q}{m\gamma v c} \vec{p} \times \vec{B} = - \frac{q\vec{B}}{m\gamma v c} \times \vec{p}$$

Thus, \vec{p} rotates with angular velocity $\vec{\omega}_{\text{cyclotron}} = - \frac{q\vec{B}}{m\gamma v c}$ and the particle moves in a circle

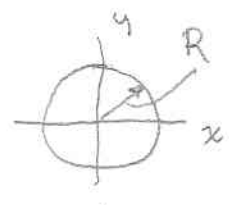
$$x(t) = R \cos \omega_c t \quad y(t) = R \sin \omega_c t$$

$$v_x = -R\omega_c \sin \omega_c t \quad v_y = R\omega_c \cos \omega_c t$$

$$v = \omega_c R \quad p = m\gamma v = m\omega_c R\gamma v$$

$$\text{or } p = m \cdot \frac{qB}{m\gamma v c} \cdot R\gamma v$$

$$R = \frac{pc}{qB}$$



5. How does a particle with charge q move in a uniform magnetic field \vec{B} ?

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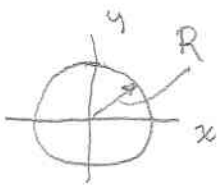
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$$R = \frac{pc}{qB}$$

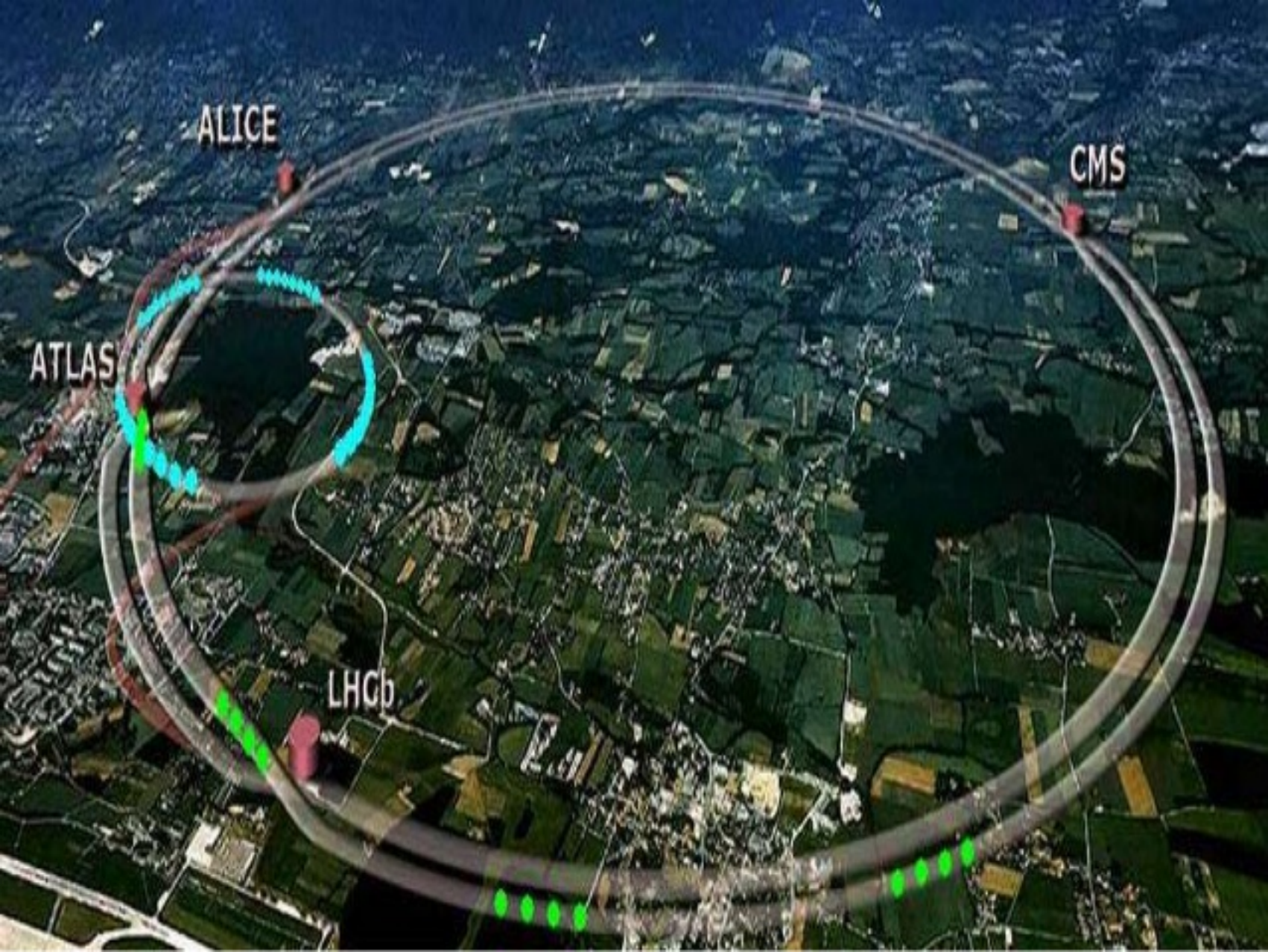


Problem: The Large Hadron Collider (LHC) at CERN uses $B = 8.36$ Tesla and accelerates protons to $7 \text{ TeV} \equiv 7 \times 10^{12} \text{ eV}$,

How large is it?

$$\begin{aligned} R &= \frac{pc}{qB} = \frac{7 \times 10^{12} \text{ e} \times \frac{1}{300} \text{ stat volt} \frac{\text{esu}}{\text{cm}}}{\text{e} \times 8.3 \times 10^4 \frac{\text{esu}}{\text{cm}^2}} \\ &= \frac{7}{3 \times 8.3} \times 10^6 \text{ cm} = 2.8 \text{ km} \end{aligned}$$

(actual radius is 4.3 km)



ALICE

CMS

ATLAS

LHCb

Problem: The Large Hadron Collider (LHC) at CERN uses $B = 8.36$ Tesla and accelerates protons to $7 \text{ TeV} \equiv 7 \times 10^{12} \text{ eV}$.

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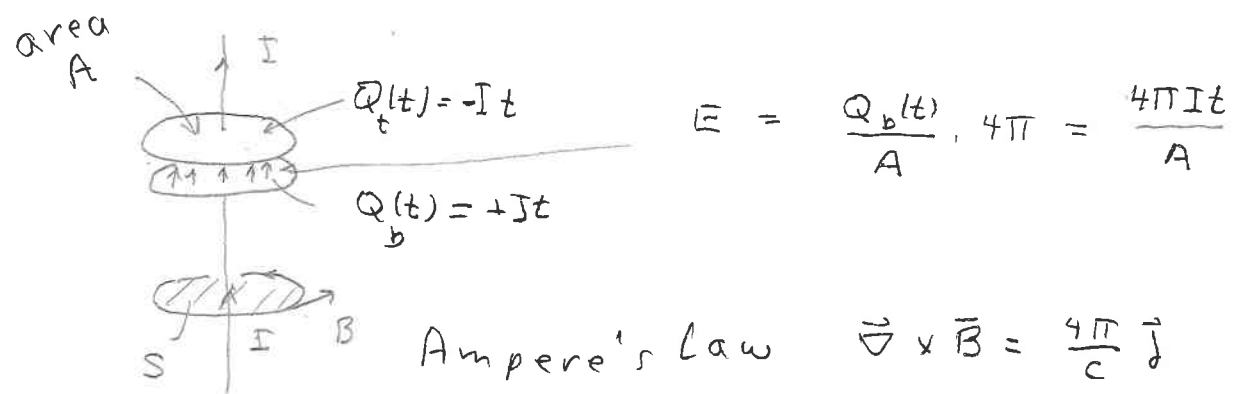
$$R = \frac{pc}{qB} = \frac{7 \times 10^{12} \text{ e} \times \frac{1}{300} \text{ stat volt}}{\text{e} \times 8.3 \times 10^4 \frac{\text{esu}}{\text{cm}^2}}$$

$$= \frac{7}{3 \times 8.3} \times 10^6 \text{ cm} = 2.8 \text{ km}$$

(actual radius is 4.3 km)

C Now consider changing fields

1. Maxwell's displacement current



$$E = \frac{Q_b(t)}{A}, 4\pi = \frac{4\pi I t}{A}$$

$$\text{Ampere's Law } \nabla \times \vec{B} = \frac{4\pi}{c} \vec{j}$$

$$\oint_{\partial S} \vec{B} \cdot d\vec{l} = \frac{4\pi}{c} \int_S \hat{n} \cdot \vec{j} dS = \frac{4\pi}{c} I$$

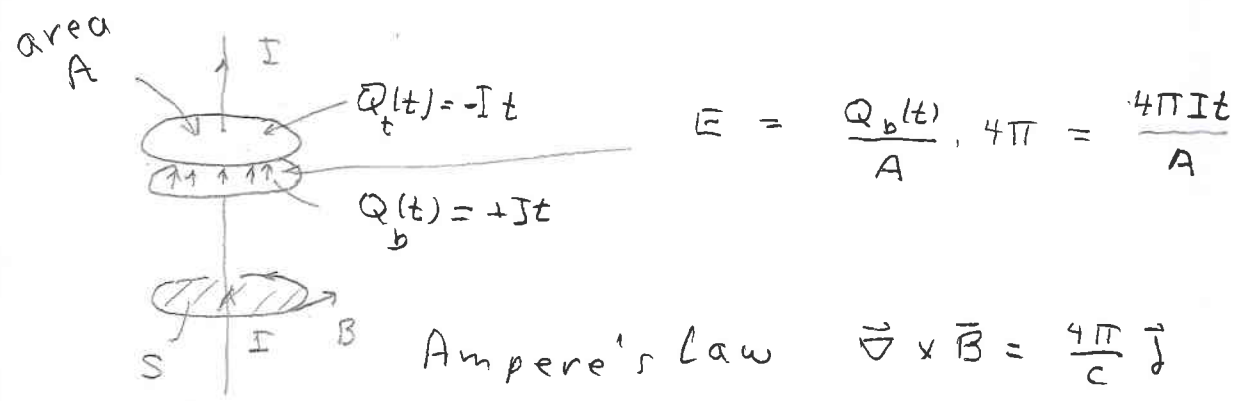
However, in this problem we could choose S to pass between capacitor plates where $\vec{j} = 0$ and find a contradiction: $\vec{B} = 0$??

$$\text{But we forgot } \nabla \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t} !$$

$$\begin{aligned} \oint_{\partial S'} \vec{B} \cdot d\vec{l} &= \frac{1}{c} \int \frac{d\vec{E}}{dt} \cdot \hat{n} dS = \frac{1}{c} \frac{d}{dt} \left(\frac{4\pi I}{A} t \right) A \\ &= \frac{4\pi}{c} I \quad \text{Problem solved} \end{aligned}$$

C Now consider changing fields

1. Maxwell's displacement current



Ampere's Law $\vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j}$

$$\oint_{\partial S} \vec{B} \cdot d\vec{l} = \frac{4\pi}{c} \int_S \hat{n} \cdot \vec{j} dS = \frac{4\pi}{c} I$$

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But we forgot $\vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$!

$$\oint_{\partial S'} \vec{B} \cdot d\vec{l} = \frac{1}{c} \int_{S'} \frac{d\vec{E}}{dt} \cdot \hat{n} dS = \frac{1}{c} \frac{d}{dt} \left(\frac{4\pi I}{A} t \right) A$$

$$= \frac{4\pi}{c} I \quad \text{Problem solved}$$

This consistency is guaranteed:

$$\vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{1}{c} \frac{\partial \vec{E}}{\partial t}$$

take the divergence of both sides

$$0 = \vec{\nabla} \cdot [\vec{\nabla} \times \vec{B}] = \frac{4\pi}{c} \vec{\nabla} \cdot \vec{j} + \frac{1}{c} \underbrace{\vec{\nabla} \cdot \frac{\partial \vec{E}}{\partial t}}_{\frac{\partial}{\partial t} \underbrace{\vec{\nabla} \cdot \vec{E}}_{4\pi \rho}}$$

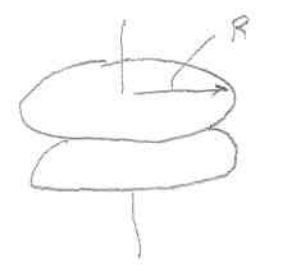
$$\vec{\nabla} \cdot \vec{j} + \frac{\partial \rho}{\partial t} = 0$$

$$= \frac{4\pi}{c} \left(\underbrace{\vec{\nabla} \cdot \vec{j} + \frac{\partial \rho}{\partial t}}_{=0} \right)$$

is required for

$$\vec{\nabla} \times \vec{B} = \frac{4\pi}{c} \vec{j} + \frac{\partial \vec{E}}{\partial t} \text{ to make sense!}$$

3. Find \vec{B} field between plates of charging capacitor



top view



\vec{B} will be tangent to circles concentric with symmetry axis

$$\oint_{\partial S} \vec{B} \cdot d\vec{r} = \frac{1}{c} \int_S \frac{\partial \vec{E}}{\partial t} \cdot \hat{n} dS$$

$$2\pi r B(r) = \frac{1}{c} \frac{I}{\pi R^2} \times 4\pi \pi r^2$$

$$B(r) = \frac{2I}{R^2} r$$